

***AN INTRODUCTION TO BASIC PHENOMENA OF PLASMA PHYSICS***

Dr. Arvind Prasad

\*Assistant Professor, \*\*Associate Professor

**ABSTRACT**

*Plasma is a set of neutral and charged particles which reveals a number of collective behaviors. The very long range coulomb forces enable the charged particles in plasma to work together with one another simultaneously. The study of plasma is actually a really ancient area of investigation in plasma physics and it remains to be among the vital fields due to the crucial role of its in most plasma uses including plasma processing, fabrication of semiconductor systems, etching, etc. except the presence of just ions and electrons, the plasma in many instances, has a number of other species of ions like negative ions which impact the complete plasma behaviour. Within this paper we study about the fundamental ideas of*

*plasma physics.*

**Keywords:** *Plasma, ions, hot, cold.*

**1. INTRODUCTION**

The basic plasma consists of mainly the ions as well as electrons. Nevertheless, a multicomponent plasma consists of several charged species i.e. extra positive ions, negative ions, electrons as well as dust particles. Dusty plasma as well as plasma with negative ions is actually 2 vital fields of multicomponent plasma at current. In the current study both the instances of multicomponent plasma is actually taken into consideration. The dusty plasma is actually a crucial department of plasma physics which covers a broad range of uses from astrophysical plasma to fusion plasma. On the flip side the damaging ion plasma plays an important role in the negative ion based neutral beam injection (NBI) devices used for plasma heating in the fusion reactors. The presence of negative ions can also be substantial for the fabrication of microelectronics of the semiconductor business.

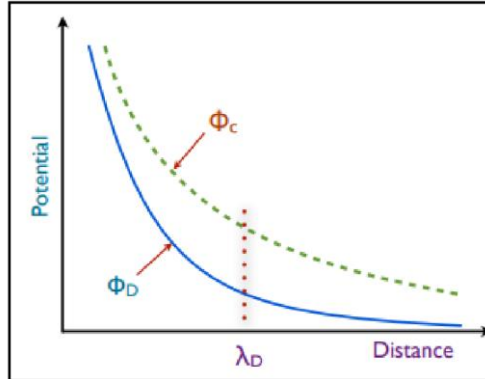
The behaviour of plasma may be studied by following 2 typical approaches: fluid and kinetic approach. Within the fluid approach, the motion of fluid components is actually taken into consideration. And fluid equations are resolved for every species. Whereas for kinetic approach, the evolution of the velocity division perform in stage room is recognized as. In certain instances both the techniques are believed to be together exactly where several species are actually handled with sleep and fluid approach kinetically according to the timescale of the phenomenon found. Occasionally neither fluid nor kinetic approach is actually appropriate to describe the behaviour of the plasma. In case which is that that it becomes crucial that you monitor the person particle trajectory that is a hard process.

**2. BASIC OF PLASMA PHYSICS**

Let us assume that a positive test charge  $q_t$  is actually positioned within the plasma at  $r = 0$ . The test charge attracts all of the negative charge and repels the positive charge. The charge density of the electron is going to increase near the test charge, even though the ion selection density will lessen. Hence, it is going to modify the charge

density distribution around an exam charge. The electrical potential around an exam particle is driven from Poisson's equation.

$$\nabla^2 \phi(r) = \frac{e}{\epsilon_0} [n_e(r) - n_i(r)] - \frac{q_t}{\epsilon_0} \delta(r - v_0 t),$$



**Figure 1: Comparison of the Debye and Coulomb potentials**

wherever  $e$  is actually the magnitude of the electron charge, the permittivity of the free room,  $n_e$ ,  $n_i$  are actually the electron as well as ion number densities, respectively,  $(r-v_0)$  the

Dirac delta feature, as well as  $v_0$  a frequent velocity of an exam charge in case it moves. For a stationary test charge, we are able to set  $v_0$  to zero.

$$= \frac{q_t}{4\pi\epsilon_0 r^2}$$

wherever  $n_0$  is actually the equilibrium electron number density,  $k_B$  the Boltzmann constant, as well as  $T_e$  the electron heat. Much from origin,

we've  $\beta q \phi \ll 1$ . Therefore, the Fourier transformation of Poisson's situation yields

$$= 8\pi n_0 q + \frac{\exp(ik \cdot r)}{k^2 D}$$

where we have denoted  $2 = 1 + 1/5\epsilon_s$ , with the electron Debye radius given by

$$5\epsilon_s = \frac{7}{k^2 D}$$

The prospective division around a test stationary ion is then the shielded Coulomb possible

$$= \frac{1}{56s} \exp - \frac{1}{r} \phi$$

The prospective division around a positive test charge is likewise recognized as the Debye Huckel or maybe the Yukawa potential.

influences are actually overlooked and a number of approximations are actually taken into consideration, a fluid description referred to as MHD (Magneto Hydro Dynamics) is actually sufficient to describe the equilibrium & balance qualities associated with favorite plasma.

### 3. TYPES OF PLASMAS

Based on the heat of the person charged species, plasmas could be split into 2 categories:

1. Hot Plasmas (or high temperature plasmas)
2. Cold Plasmas (or low temperature plasmas)

#### *Hot Plasmas*

The high temperature plasmas are actually plasmas which are fully ionized, have very limited amount and few collisions of kinetic strain. For example: Astrophysical plasmas like stars and magnetically confined fusion plasmas. The distribution feature might deviate considerably from Maxwellian because of the low collision frequency as well as extended mean open path of the particles. Thus it's becomes important to explain the high temperature plasmas with a kinetic therapy. In case the above mentioned



Figure2: Examples of high temperature plasma: the fusion plasma in the tokamak ASDEX (T = 10keV)

*Cold Plasmas*

The low temperature plasmas are actually split into non-thermal and thermal plasmas. The plasma in which the particles are actually at exact same temperature or perhaps at local thermodynamic equilibrium (LTE) is actually recognized as winter plasma. For example: arc discharge at atmospheric pressure. The standard heat of a winter plasma is,  $T_e \approx T_{ion} \approx T_{gas} \approx 10^4K$

**CLASSICAL AND QUANTUM PLASMAS**

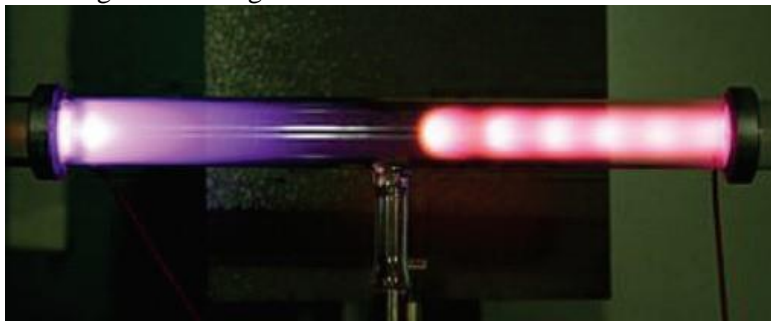
Plasma physics has huge uses in low temperature laboratory plasmas (gas discharges), in higher power density plasmas, e.g. for example Inertial Confinement Fusion (ICF) schemes, in Magnetic Confinement Fusion (MCF) schemes, in magneto inertial fusion schemes, in our solar energy system (viz. geospace plasmas, heliospheric plasmas), in planetary methods, and in



**Figure3: Arc discharge of Ar with H<sub>2</sub> plasma (Thermal plasma).**

In a non thermal plasma, the electron temperature is significantly greater compared to the heat of the ions as well as neutrals as well as the amount of ionization is comparatively small ( $10^6$   $10^4$ ). A low pressure glow discharge is a

good example of non thermal plasma in which the standard heat is,  $T_e \approx 10^4K \gg T_{ion} \approx T_{gas} \approx 300K$  (Partial thermodynamic equilibrium, PTE).



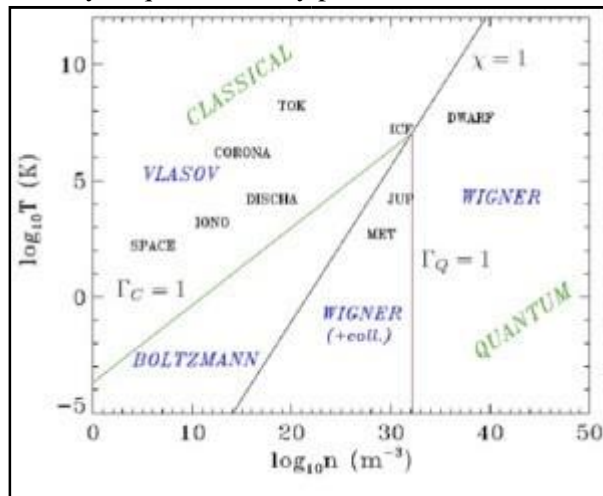
**Figure4: Glow discharge tube (Non-thermal plasma)**

**4. PHYSICAL REGIMES FOR astrophysical environments (e.g. interstellar**

medium, the cores of white dwarf stars, etc.), neutron stars, magnetars, and within MHD energy transformation. The plasma fusion efforts will provide a cost effective, limitless and environmentally friendly alternative energy, because the method would not produce unwanted waste fusion products, and in most cases would expose insignificant hazards to mankind. Plasmas can also be used for medical therapies as well as for enhancing adhesion, printing, and lamination.

Classical plasma physics is primarily centered on regimes of low densities in addition to higher temperatures. More than numerous years, new emerging subfields of plasma physics have been growing quickly, such as clearly coupled dusty and quantum plasmas, ultra cold basic plasmas, and intense laser plasma/solid density plasma interactions. In the following part, we'll briefly

characterize particular bodily parameters for investigation of ours. In paper six, we've quantumdusty plasmas as well as the programs examined an innovative strictly raising ion of theirs, which have been a concern in the streaming instability in quantum dusty plasma.



**Figure 5: Plasma diagram in the  $\log T - \log n_0$  plane, separating the quantum and classical regimes. METAL: electrons in a metal; IONO: ionospheric plasma; TOK: plasma in the typical tokamak experiments for nuclear fusion; ICF: inertial confinement fusion; SPACE: interstellar plasma; DWARF: white dwarf star.**

### 5. BASIC PARAMETERS FOR QUANTUM PLASMAS

The quantum appearance in plasmas start to be essential if the de Broglie length  $\lambda_D$  is actually akin to the Thomas Fermi radius  $r_{TF} \leq \lambda_D$ .

The distinctive de Broglie length in a thick plasma with degenerate electrons is actually provided by  $\lambda_D = \hbar / (k_B T_e)$ , in which  $k_B$  is

$$\lambda_D = \left( \frac{\hbar^2}{2k_B T_e m_e} \right)^{1/2} = \left( \frac{\hbar^2}{2m_e k_B T_e} \right)^{1/2}$$

$\hbar < \lambda_D$

$$\Gamma_E = F \frac{C}{H} = F \frac{1}{H}$$

which identifies the ratio between the plasmonic energy related with the electron plasma oscillations and also the Fermi electron energy. The distinctive de Broglie length could be conveyed in phrases of the quantum coupling

parameter as  $\Gamma_E = \lambda_D / r_{TF}$ . For  $\Gamma_E \geq 1$ , the quantum consequences start to be essential,

result, classical and quantum routines take place in various physical environments.

### 6. CONCLUSION

Plasma is present in several forms in nature and features a prevalent use of technology and science. Owing to the presence of free ions, utilizing plasma for ion resources is pretty common. Because of this specific case, plasma is made by a good type of low pressure gasoline discharge. This particular paper has a brief discussion of a few essential plasma phenomena.

Fermi electron velocity,  $v_F = C$ , is actually the Fermi electron energy, as well as  $\sim \hbar / \lambda_D$  the Plank constant divided by  $\lambda_D$ . Of course, the quantum consequences play a crucial role every time the plasma temperature  $T_e$  is actually higher compared to  $T_{Fe}$ , whereas a classical plasma description is usually sufficient in the complete opposite limit. The coupling parameter of quantum plasmas reads

representing that  $\Gamma_E$  is actually equal to or even greater than the common inter electron distance  $\lambda_D = r_{TF}$ . Nevertheless, when  $\Gamma_E \rightarrow$

$\Gamma_E$ , the quantum coupling parameter tends to the classical coupling parameter, i.e.  $\Gamma_E \rightarrow \Gamma_L$ . In the classical regime,  $\hbar \rightarrow 0$ , as well as the electrons may be seen as point like and no quantum interference outcome shows up. As a

### REFERENCES: -

[1] Adhikari, B.R. & Basnet, Suresh & Lamichhane, H.P. & Khanal, Raju. (2020). Variation of Velocity of Ions in a Magnetized Plasma Sheath for Different Magnetic Field. Journal of Nepal Physical Society. 6. 25-29. 10.3126/jnphysoc.v6i1.30513.

[2] Patel, Soniya & Varma, P. & Tiwari, M.. (2019). Plasma Physics and Controlled Fusion. 10.13140/RG.2.2.21070.10564.

- [3] Gates, David. (2018). Plasma: An International Open Access Journal for All of Plasma Science. Plasma. 1. 4. 10.3390/plasma1010004.
- [4] Conde, Luis. (2018). An Introduction to Plasma Physics and its Space Applications, Volume 1: Fundamentals and elementary processes. 10.1088/2053-2571/aae132.
- [5] Kovtun, Yuri &Skibenko, A. &Skibenko, E. &Larin, Yu. (2010). Investigation of multicomponent plasma parameters by microwave methods. 2010 International Kharkov Symposium on Physics and Engineering of Microwaves, Millimeter and Submillimeter Waves (MSMW). 10.1109/MSMW.2010.5546081.
- [6] Schollmeier, Marius & Roth, Markus &Schaumann, G &Blazevic, Abel &Flippo, Kirk & Frank, A & Fernandez, Juan & Gautier, Donald & Michel, Knut &Heßling, Thomas &Hegelich, Bjorn &Nürnberg, Frank &Pelka, A &Ruhl, Hartmut & Schreiber, Jörg & Schumacher, Dennis & Witte, K &Zielbauer, B & Hoffmann, Dieter. (2008). Plasma physics experiments at GSI. Journal of Physics: Conference Series. 112. 042068. 10.1088/17426596/112/4/042068.
- [7] Browning, Philippa. (2005). Introduction to Plasma Physics: With Space and Laboratory Applications. Plasma Physics and Controlled Fusion. 47. 1109. 10.1088/0741-3335/47/7/B01.