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Quantitative Method of Singularly Perturbed

Delay Differential Equations

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Abstract – **In this paper, singularly perturbed delay differential equation is solved by converting to first order ordinary differential equation and using two- point Gaussian quadrature rule to approximate the solution. Numerically the values of the derivatives and the integrals are approximated to get tridiagonal system. Thomas algorithm is used to solve the tridiagonal system of the method. Convergence of the proposed method is analyzed**

Keywords: **singularly perturbed delay differential equation, Taylor series, Boundary Layer,**

Tridiagonal system

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I. Introduction

Singularly perturbed differential-difference equations (SPDDEs) arise very frequently in the mathematical modeling of real life situations in science and engineering variational problems in control theory $[2]$, $[6]$, $[14]$. In general, any ordinary differential equation in which the highest derivative is multiplied by a small positive parameter and containing at least one delay/advance parameter is known as singularly perturbed differential difference equation. Lange and Miura [14]-[18] published a series of papers extending the method of matched asymptotic expansions initially developed for ordinary differential equations to obtain approximate solution of singularly perturbed differentialdifference equations. Numerical analysis of singularly perturbed differential - difference turning point problems was initiated by Kadalbajoo and Sharma. In a series of papers, see $[8]-[10]$, they gave many robust numerical techniques for the solution of such type of problems. Kadalbajoo and Sharma [8] elucidate a numerical method to solve boundary value problems for singularly perturbed differential difference equation of the form: , under the interval and boundary conditions on and on . Kadalbajoo and Sharma [9] proposed a numerical method to solve boundary-value problems for a singularly perturbed differential-difference equation of a mixed type, i.e. which contains both types of terms having negative shifts as well as positive shifts, and considered the case in which the solution of the problem exhibits rapid oscillations. Kadalbajoo and Sharma [10] described a numerical approach based on finite difference method to solve a mathematical model arising from a model of neuronal variability. Patidar and Sharma [20] combined fitted-operator methods with Micken's non-standard finite difference techniques for the numerical approximations of singularly perturbed linear delay differential equations. Kadalbajoo et al. [11] derived *ϵ*uniformly convergent fitted methods for the solution of singularly perturbed differential difference equation (SPDDE). Kumar and Sharma [13] presented a numerical scheme based on B-spline collocation to approximate the solution of boundary value problems for singularly perturbed differential-difference equations with delay as well as advance. Amiraliyev and Cimen [5] derived a numerical method for singularly perturbed boundary value problem for a linear second order delay

differential equation with a large delay in the reaction term. The authors constructed an exponentially fitted differential scheme on a uniform mesh accomplished by the method of integral identities with the use of exponential basis functions and interpolating quadrature rules with weight and the remainder term in the integral form.

II. Description of the problem

Consider a linear singularly perturbed two -point boundary value problem of the form: $\Box y$ *x* \Box \Box $\Box a$ *x* $y \times \Box \Box \Box \Box Z \Box \Box b \times y \times \Box \Box \Box \Box f \times \Box \Box, 0 \Box \times \Box 1, (1)$ on (0, 1), under the boundary conditions $y x() \Box Z(x)$ on $-\sqrt{2}x \Box 0$,

 (2) $y(1)$ \Box D ,

where \overline{D} is small parameter, $0 \Box \Box \Box \overline{D}$ 1 and \overline{D} is also delay

parameter $0 \Box \Box \Box 1$ a(),b(),x $x = f x$ (),and () $\Box x$ are bounded continuous functions in $(0,1)$ and \overline{D} is a finite constant. For \overline{D} 0 the solution of the boundary value

problem (1) and (2) exhibits layer or oscillatory behavior depending on the sign of $\Box a$ *x*() $\Box b$ *x*() . \Box If

 $\Box a \times b \times () \Box () \Box \Box \Box M \Box 0$, where *M* is a positive constant, the solution of the problem (1) and (2) exhibits layer behavior and if $\Box a$ $x()$ $\Box b$ $x()$ $\Box \Box M$ 0, it exhibits oscillatory behavior. The boundary value Problem considered here is of the reaction-diffusion type, therefore if the solution exhibits layer behavior, there will be two boundary layers which will be at both end point i.e., at *x* \Box 0 and *x* \Box 1.

III. Numerical Method

We divide the interval $0,1\Box$ \Box into *N* equal parts of mesh size h, with $0 \Box^{x} x_0$, ..., $x_{N\Box 1}$, $x_N \Box 1$ being the mesh points. Hence, we \perp

have $\binom{x}{i}$ \Box *ih*, for *i* \Box 0 *to N* . We choose n such that $\binom{x}{n}$ \Box 2. In the

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 \Box 1 \Box

Interval \Box D \Box 2 \Box the boundary layer will be at the left

 $\Box 1$

hand side i.e., at $x \Box$ 0 and in $\Box \Box_{2}$ ¹ \Box the boundary

layer will be at the right hand side i.e., at $x \Box 1$. Therefore we calculate the solution for three Different

 \Box $1 \square 1$ $\square 1$ \square intervals \Box \Box 0, $_2$ $\overline{v_0}$, $_2$ and $\overline{v_0}$, $\overline{v_0}$. *III.1 Left – End Layer* For left hand layer $\mathbf{1}$ i.e., $0 \Box \Box x$ 2, by using Taylor's expansion for *y x* $\Box \Box$ $\Box \Box$, we can

say that, $y \Box x - \Box \Box y \times \Box \Box \Box x$ Hence,

 Dy 00x $D \cup y$ *x* $D \cup y$ *x* $D \cup y$ *x* $D \cup D$

By using this approximation in (1) we get,

 y *x* \Box \Box \Box y *x* \Box \Box \Box \Box a *x* y *x* \Box \Box a \Box b *x* y *x* \Box \Box f *x* \Box \Box f *x* \Box \Box g *yx* \Box a b *x* b *yx* \Box a b *x* b *yx* \Box b *yx* \Box b *yx* \Box b *yx* \Box b *yx* (3)

Inregrating this equation from *x to x* $i \square_1$ we get, $x_i \Box 1$ $x_i \Box 1$ $x_i \Box 1$ $x_i \Box 1$

 y_i _{*u*} \Box *y*_{*x*} \Box *D* \Box *<i>D* \Box *<i>dx* \Box *D* \Box *ax y x* \Box *D* \Box \Box *f x dx* \Box \Box

xi xi xi xi (4)

We use Gause-Legendre quadrature to approximate the value of this integration, for which we need to change the limits to -1 to 1. For this purpose, we introduce a new variable u Where,

2 - $x \Box x_i \Box x_i_{i \Box_1} \Box u h \Box \Box x_i \Box x_{i \Box_1} \Box u \Box$ $\Box x \overline{\Box x}i\overline{\Box 1 \Box x}i$ 2 Using this approximation in (4) we get $y_i \Box_1 \Box y_i \Box$ *y*(x_i ^{\Box *D*) \Box *y*(x_i \Box *D*) (5)} Using two- point Gaussian quadrature we get

1 0100010 *f x dx f f* , $\overline{1}$ $\overline{1}$ $\sqrt{3}$ $\overline{1}$ $\sqrt{3}$ $\overline{1}$ $\sqrt{3}$ $\overline{1}$

For any function f, continuous and differentiable in interval [- 1,1].

when $u \Box \frac{1}{\Box}$, 3 $\mu \Box \neq 0$

$$
\frac{ih \Box x_i \Box x_{i\Box L}}{a \Sigma_{i\Box L} \Box x_{i\Box L}} \frac{h \Box \Box \Box \Box}{2} \Box x_{i\Box L} \Box k,
$$
\nwhere $k \Box \frac{h \Box \Box \Box \Box}{2}$,

$$
\frac{\Box 1}{\sqrt{3}}, \frac{uh \Box x_i \Box x_{i\Box 1}}{2} \Box x_i \Box^{-1} \frac{h \Box \Box \Box}{2} u \Box x_i \Box k,
$$
\n
$$
\frac{\Box 1 \Box}{h \Box 1 \Box \Box}
$$
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h \Box 1 \Box \Box
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$$
h \Box 1 \Box \Box
$$
\n
$$
\frac{\Box 1 \Box}{2}
$$
\nWhere k \Box \qquad \frac{\Box}{2}

Substituting this in (5) we get y_i ⁿ₁ \Box $\Box y_i y x \Box_{i \Box 1} \Box \Box \Box x y x \Box_i \Box \Box$ $h \Box a$ $x \Box$ $i \Box k$ y $x \Box$ \Box $i \Box$ \Box $\Box k$ $\Box a$ $x \Box$ $i \Box k$ y $x \Box$ \Box $i \Box k$ $\Box a$ \Box \Box 2 *b x ⁱ* ¹ *k y x ⁱ* ¹ *k b x ⁱk y x ⁱk f x ⁱ* ¹ *k* $f x \Box_i \Box k \Box \Box_0$

(6)

Now we use following approximations for the y values, $y \times \Box$ i_{01} \Box k \Box \Box \Box \Box y_i \Box \Box y_i \Box \Box y_i \Box \Box y_i \Box \Box y_i

k yⁱ ¹*yⁱ y* , *yi* 1 *yⁱ* ¹. *h h i* 1 *h k k*

Similarly,

y x i k 1 *k h yi k h yi* 1,

k k y ⁱ 1 *y x i* 1 *k yi* 1, *h h ky x i k* 1 *yi k yi* 1 *h h* By substituting this in (6) we get, *h f x ⁱ* ¹ *k f x ⁱk* 2 *k k h a x i k* 2 *b x i k* 2 *yi* 1 *x i* 1 *k* 2 *b x k i* 2 *yi*

 \Box $\Box \bar{\mathcal{Q}}$ $h k\Box \Box D$ $h k\Box$ $\Pi^0 \Pi_h \Pi a \; x \Pi_{i01} \Pi k^{\Pi}$ $_2$ $\Pi b \; x \Pi_{i01} \Pi k^{\Pi} {}_2{}^{\Pi} {}_0 y_{i01}$. (7)

2 *k h k k h k h a x i* 1 *k* 2 *a x i k* 2 *b*

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We can say that,

 $A_i y_i$ _{\Box} $B_i y_i \Box C_i y_i$ _{\Box} b_i , For $1 \square \square i$ *n* $\square 1$, where \varPi $k\square\varPi$ *k* $A_i \Box \Box a \times k \Box_i \Box \Box$ $\Box b \times k \Box_i \Box \Box h$, 2 $\overline{2}$ $h \Box k \Box T$ $\Box 2 \Box$ $k \Box \Box$ $B_i \Box$ \Box \Box $a x \Box$ $i \Box$ $k \Box$ h 2 2 k *h* \Box *k b x ⁱ* ¹*k b x ⁱk* , 2 2 D *h k* **0** *h k* **0** *c_i* **0** *a**x***0** *j*₀₁*a**k***0** $\Box b \times \Box_{i} \Box k \Box$ $\Box h \quad 2$ 2 *h f xi* 1 *k f xi k b ⁱ*. τ

$$
\overline{1}
$$

 \mathbf{I} i.e., 2DxD1, by using Taylor's expansion for $y \Box \Box x \Box \Box \Box x$ we get, $y \Box \Box x$ $\Box \Box \Box \Box$ $y \Box \Box x \Box \Box \Box y \Box \Box \Box x \Box$ Hence, *<i><i>z*_y 00x 00 *y* 02x 02x *y* 00x 0 By using this approximation in (1) we get, *y x* \Box \Box \Box *y x* \Box \Box $\Box Z$ \Box \Box $\Box Z$ $\Box Y$ $\Box Y$

III.2 Right – End layer For right hand layer

Integrating (8) from x_i ^{α}*i* to x_i gives us,

xi xi xi xi

yi yⁱ ¹ *y x dx a x y x dx b x y x dx* \Box \Box *f x dx* \Box \Box *xi* 1 *xi* 1 *xi* 1 *xi* 1

$$
(9)
$$

Here the value of u for using Gauss-Legendre Quadrature will be,

(10) Now we use following approximations for the y values y $x \Box$ $\Box k$ $\Box \Box \Box$ \Box y_i $\Box \Box k$ $\Box \Box \Box$ y_i \Box ,

$$
\Box y_i \Box, \quad \frac{\Box k \Box Z \Box \Box y_i \Box y_i \Box}{h}
$$
\n
$$
\Box \quad \frac{k \Box Z \Box}{h} \quad \frac{k \Box Z}{\Box \Box y_i \Box h} \quad y_i \Box.
$$

Similarly,

y x i-1 *k* 1 *k k yi*, *y x ⁱk* 1 *yⁱ* ¹ *h h* ¹, *ⁱ ⁱ k k y y h h k*

 y x i 1 *k* 1 *yi* 1 *k yi*. \Box *h* \Box *h* By substituting this in (10) we get, $hf x k \Box \Box I$ $_i \Box$ $f x \Box$ $_{i \Box 1} \Box k \Box \Box \Box D$ *h k* \Box $i \Box$ $h k \Box$ \Box \square \square 2 **h** $a x\Box i\Box 1 \Box k\Box 2 \Box b x\Box i\Box 1 \Box k\Box 2 \Box 0 y i\Box 1$

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 $h kD D$ *h* kD $h kD$ **h** kD **h** kD **k** $En \Box 2 ___2 \Box a_n \Box b_n$ *h h a x k i* 2 *a x i* 1 *k* 2 *b x k i* 2 *b x i* 1 *k* $2 \Box \Box$ *yi* \varOmega \Box *Fn* \Box 2 \Box *a_n* \overline{h} $2h\overline{b}n \Box fm$. \Box $k\Box Z$ $k\Box$ *N* =20,40,80,160,and 320,for different values of *h a x k i* 2 *b x k i* 2 *yi* 1 \Box and \Box \Box TABLE I We can say that, THE MAXIMUM ERRORS IN SOLUTION OF $G y_{i i} \Box H y_{i i} \Box I y_{i i} \Box b_i$, for Π EXAMPLE 1 WITH \Box 0.1 i N-1, where , $M = 100$ $N = 200$ $N = 300$ $N = 400$ $D(11)$ *h k* D D *h k* D __ 0.01 3.157 e-003 1.605e-003 1.075e-003 8.083e-010 G_i **d** \Box **a** x **D** i **n**₁**D** k **D** \Box **b** x **D** i **D** i **D** i 0.04 3.143 e-003 1.599e-003 1.075e-003 8.033e-004 $R = \frac{h}{2}$ 2 2 0.06 \Box 2 \Box a x \Box i \Box 1 \Box 3.078 e- \Box 2*Dh* k \Box *DHk* \Box *Dhk*Finally we can write the formulat \Box *k* ion in the matrix form *Hi* \Box *m* \Box *a* \Box *i* \Box *k* \Box *k* 2 *b x i k* as, 2 *Ay b x bi*, 1where *k* 2 003 1.566e- $DB1$ C_1 $0 \qquad \cdots$ \cdots \cdots $0\Box$ 1 003 $Dh k$ $D Dh k$ I_i $D = D a x$ I_i D_i Dk D_i D_i D_i Π \square 2 医细胞 医手 B_2 \cdot C_2 $\mathcal{L}_{\mathcal{L}}$ 1.050e $x \Box$ $i \Box$ *k* \square $h \ 2 \ 2 \ h \ f \ x \Box$ \square $h \ \Box$ $f \ x \Box$ $i \Box$ $h \ \Box$ $h \ \Box$ A \Box \mathcal{L}_{max} , and \mathcal{L}_{max} 003 \Box 2 *An*□1 *Bn*□1 *Cn*□1 7.9001e-2 in a shekarar 1990 \Box 0 014 \Box \Box \Box : *g a x a x a n x a n a x a n d <i>x a n b x n n g x n n g x n n y x n n n y x n n n y x n n n y x n n n y x n n n y Hn* 1**i** 1**In** 1**i** 1 \Box \Box (12) GN-2 HN-2 IN-2 We use following approximations in this equation. $y \square y$ \Box \Box *y x ⁿy x ⁿy x ⁿyⁿ n* 1 *n* 1 , 2*h* \Box 0 0 *GN* \Box 1 *HN* \Box 1 *y* **□**2*y* **□** *y y*[□] ^D*y*¹ ···*yN*□1□□, *y***u**_n 00 ___________________*n*⁰1 ₂*n* \Box $n \Box 1$ *h b* \Box *Db*1 \Box *A*1*y*0 **b**2 ⋅⋅⋅ **bN-2 bN-1** \Box *CN* \Box 1*yN* \Box To get, \Box *y* \Box 2*y* \Box *y* **IV. Numerical Example** *fn n* 1 2*nn* 1 *an yn yn* 12 *hyn* 1 *b yn n* An example of this type of equation is, \Box \Box *h* \Box Δy \Box $\Box x$ $\Box 2y$ $x \Box$ $\Box x \Box y$ $x \Box$ \Box 1, $y \Box$ \Box 1, $y \Box$ \Box 1, $y \Box$ Rearranging this equation gives us, (13) *A D E F* , For $x_n \Box$, $n \quad n \quad n$ \Box \Box \Box \Box \Box $D \Box$ $D \Box Z \Box x \Box 0$ and $y 1 \Box \Box \Box 1$ *fn h*2 *an* 2*h yn* 1 2 *h*2 *an bn yn h*2 *an* 2*h yn* For 1this equation, solution was calculated for We can say that, (14) *Dn* $\forall n \Box 1 \Box E_n$ $\forall n \Box F_n$ $\forall n \Box 1 \Box bn$, Where, *Dn* 2 *an* , *h* $\sqrt{ }$ 2*h* $\mathcal{D}% =\mathcal{L}_{1}\cup\mathcal{L}_{2}$

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Fig.1.Numerical solution for π 0.1 and π 0.01

Fig.2.Numerical solution for $\pi/2$ 0.01 and $\pi/2$ 0.04

Where u is the exact solution, u_h is the numerical solution is a constant, h is the mesh size and *p* is the order of convergence.

Here, as the exact solution is not knows we assume the numerical solution at $N = 320$ to be the exact solution and use other numerical solutions to plot log $u \mu \Box$ *h* versuslogh, for

 \Box 0.01 and \Box 0.04 ,to obtain the convergence plot. Slope of convergence plot is equal to the order of convergence.

From this plot, it can be said that order of convergence for this method is 1.

V. Conclusion

The numerical analysis of singularly perturbed delay differential equation is discussed. In general numerical solution of second order differential equation will be more difficult than numerical solution of first order differential equation. Even though the numerical results are computed at all points of the mesh size *h* , only few results have been reported. It is observed that the present method approximates the exact solution very well for $h\square\mathcal{L}$ for which other classical finite difference methods fail to give good results. The effect of the delay parameter on the solution has also investigated and presented by using graphs. Convergence of the proposed method is analyzed. The present method is simple, very easy to implement and efficient technique for solving singularly perturbed delay differential equations.

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